# Resistivity of a sinusoidally corrugated surface 

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## Introduction

From ChrisA's talk. Paper by S. Morgan, Journal of Applied Physics 20, 352 (1949).

## Effect of Copper Roughness on Resistivity



There is an important parameter missing in this plot: corrugation slope $\sim \Delta / \lambda$.

## Model

Assume the sin profile of the surface

$$
y_{0}(x)=h \cos \left(\frac{2 \pi x}{\lambda}\right)
$$

These are grooves parallel to $z$-axis.
5 roughness profiles


For $2 \pi h / \lambda=4$ the steepest slope angle is $76^{\circ}$.

## Direction of the magnetic field

Magnetic field can be considered uniform at distance from the surface $\gg$ corrugation scale (but much smaller than the RF wavelength). There are two independent directions (polarizations) of the magnetic field. Far from the surface the magnetic field approaches a constant value

$$
\boldsymbol{H}=\hat{x} H_{0}
$$

or

$$
\boldsymbol{H}=\hat{z} H_{0}
$$

x-polarization

z-polarization


## Small skin depth limit

Assume small skin depth, $\delta \ll h, \lambda / 2 \pi$. The boundary condition for the magnetic field on the surface is

$$
\left.H_{n}\right|_{\text {surface }}=0
$$

Resistivity increase factor

$$
\eta=\frac{\int_{\text {period }} H_{t}^{2} d s}{H_{0}^{2} \lambda}
$$

We will calculate $\eta$ for both polarizations, $\eta_{x}, \eta_{z}$. If magnetic field is randomly oriented relative to the grooves, averaging over all possible orientations gives

$$
\eta=\frac{1}{2}\left(\eta_{x}+\eta_{y}\right)
$$

For the $z$-polarization the magnetic field penetrates the grooves, $\boldsymbol{H}(x, y)=\hat{z} H_{0}$ for $y>y_{0}(x)$, and

$$
\eta=\frac{1}{\lambda} \int_{\text {period }} d s
$$

## Small skin depth limit $-x$ polarization

For $x$-polarization one has to solve the Poisson equation above the metal, $y>y_{0}(x)$ :

$$
\boldsymbol{H}(x, y)=-\hat{\mathbf{z}} \times \nabla \psi(x, y), \quad \Delta \psi=0
$$

with the boundary conditions

$$
\psi_{\text {surf }}=\text { const },\left.\quad \psi\right|_{y \rightarrow \infty} \rightarrow H_{0} y
$$

The function $\psi$ is periodic along $x$ with the period $\lambda$. I used computer code FreeFem++ (http://www.freefem.org/ff++/index.htm) to numerically solve the Poisson equation.

## Case $2 \pi h / \lambda=2$



Field lines


## Magnetic field on the surface



## Increase of resistivity



## How to solve 3D case in the limit of small skin depth?

3D geometry metal surface

$$
y=h(x, z)
$$

Magnetic field in free space

$$
\boldsymbol{H}(x, y, z)=\nabla \psi(x, y, z), \quad \Delta \psi=0
$$

Boundary condition on the surface

$$
\left.\frac{\partial \psi}{\partial n}\right|_{y=h(x, z)}=0
$$

Boundary condition at infinity, $\boldsymbol{H} \rightarrow \hat{\boldsymbol{x}} \mathrm{H}_{0}$

$$
\left.\psi\right|_{y \rightarrow \infty} \rightarrow H_{0} x
$$

Solve the Poisson equation in a large $0<x<a, 0<z<b$ area (say, with periodic boundary conditions) and compute $\int_{\text {surf }} H_{t}^{2} d S / H_{0}^{2} a b$.

## Finite skin depth case, z-polarization

For the $z$-polarization the magnetic field penetrates the grooves, $\boldsymbol{H}(x, y)=\hat{\mathbf{z}} H_{0}$ for $y>y_{0}(x)$. In the metal

$$
\Delta H_{z}=\frac{2 i}{\delta^{2}} H_{z}
$$

From Morgan's paper

$$
\eta=\frac{2}{H_{0} \lambda \delta} \operatorname{Im} \int_{\text {period }} d x \int_{-\infty}^{y_{0}(x)} H_{z} d y
$$

## Finite skin depth case, z-polarization

Case $2 \pi h / \lambda=2,2 \pi \delta / \lambda=0.6$

Contour lines of $\mathrm{Re} \mathrm{H}_{\mathrm{z}}$


Contour lines of $\operatorname{Im~} \mathrm{H}_{\mathrm{z}}$


## Finite skin depth case, z-polarization

This is the case studied in Morgan's paper (for different profiles of grooves).


Note that $\lambda / 2 \pi \delta=h / 2 \delta$. In the limit $\delta \rightarrow 0$ we previously found $\eta_{z}=1.68$.

## Finite skin depth case, $x$-polarization

Take

$$
\boldsymbol{H}=-\hat{\mathbf{z}} \times \nabla \psi(x, y)
$$

At $y \rightarrow \infty$ we have $\psi \rightarrow H_{0} y$, and at $y \rightarrow-\infty$ we have $\psi \rightarrow 0$. The equation for $\psi$

$$
\Delta \psi=s \frac{2 i}{\delta^{2}} \psi
$$

where the function $s=1$ in the metal and zero otherwise.

## Finite skin depth case, x-polarization



In the limit $\delta \rightarrow 0$ we previously found $\eta_{x}=1.27$

