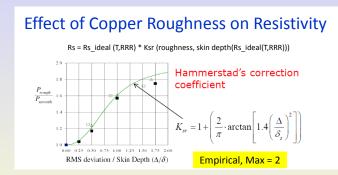
# Resistivity of a sinusoidally corrugated surface

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From ChrisA's talk. Paper by S. Morgan, Journal of Applied Physics 20, 352 (1949).



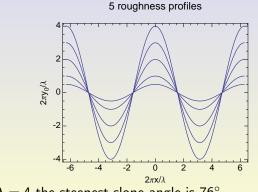
There is an important parameter missing in this plot: corrugation slope  $\sim \Delta/\lambda$ .

#### Model

Assume the sin profile of the surface

$$y_0(x) = h \cos\left(\frac{2\pi x}{\lambda}\right)$$

These are grooves parallel to z-axis.



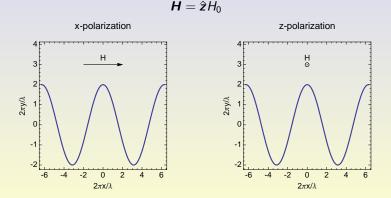
For  $2\pi h/\lambda = 4$  the steepest slope angle is 76°.

# Direction of the magnetic field

Magnetic field can be considered uniform at distance from the surface  $\gg$  corrugation scale (but much smaller than the RF wavelength). There are two independent directions (polarizations) of the magnetic field. Far from the surface the magnetic field approaches a constant value

$$\boldsymbol{H} = \hat{\boldsymbol{x}} H_0$$

or



## Small skin depth limit

Assume small skin depth,  $\delta \ll h, \lambda/2\pi$ . The boundary condition for the magnetic field on the surface is

 $H_n|_{\text{surface}} = 0$ 

Resistivity increase factor

$$\eta = \frac{\int_{\text{period}} H_t^2 \, ds}{H_0^2 \lambda}$$

We will calculate  $\eta$  for both polarizations,  $\eta_x$ ,  $\eta_z$ . If magnetic field is randomly oriented relative to the grooves, averaging over all possible orientations gives

$$\eta = \frac{1}{2}(\eta_x + \eta_y)$$

For the z-polarization the magnetic field penetrates the grooves,  $H(x, y) = \hat{z}H_0$  for  $y > y_0(x)$ , and

$$\eta = \frac{1}{\lambda} \int_{\mathrm{period}} \, \textit{ds}$$

For x-polarization one has to solve the Poisson equation above the metal,  $y > y_0(x)$ :

$$\boldsymbol{H}(x,y) = -\hat{\boldsymbol{z}} \times \nabla \psi(x,y), \qquad \Delta \psi = 0$$

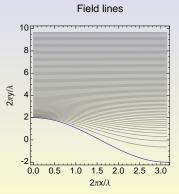
with the boundary conditions

$$\psi_{\text{surf}} = \text{const}, \qquad \psi|_{y \to \infty} \to H_0 y$$

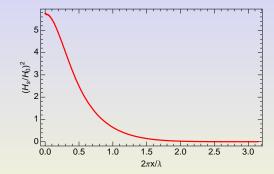
The function  $\psi$  is periodic along x with the period  $\lambda$ . I used computer code FreeFem++ (http://www.freefem.org/ff++/index.htm) to numerically solve the Poisson equation.

# Case $2\pi h/\lambda = 2$

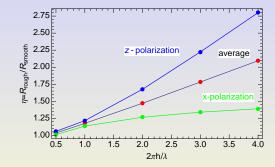




### Magnetic field on the surface



### Increase of resistivity



# How to solve 3D case in the limit of small skin depth?

3D geometry metal surface

$$y = h(x, z)$$

Magnetic field in free space

$$\boldsymbol{H}(x,y,z) = \nabla \psi(x,y,z), \qquad \Delta \psi = 0$$

Boundary condition on the surface

$$\left.\frac{\partial \psi}{\partial n}\right|_{y=h(x,z)}=0$$

Boundary condition at infinity,  $oldsymbol{H} 
ightarrow \hat{oldsymbol{x}} H_0$ 

$$\psi|_{y\to\infty} \to H_0 x$$

Solve the Poisson equation in a large 0 < x < a, 0 < z < b area (say, with periodic boundary conditions) and compute  $\int_{surf} H_t^2 dS/H_0^2 ab$ .

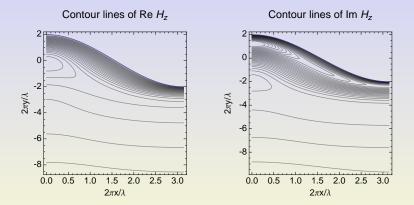
For the z-polarization the magnetic field penetrates the grooves,  $H(x, y) = \hat{z}H_0$  for  $y > y_0(x)$ . In the metal

$$\Delta H_z = \frac{2i}{\delta^2} H_z$$

From Morgan's paper

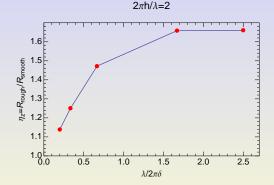
$$\eta = \frac{2}{H_0\lambda\delta} \mathrm{Im} \, \int_{\mathrm{period}} dx \int_{-\infty}^{y_0(x)} H_z dy$$

Case  $2\pi h/\lambda = 2$ ,  $2\pi\delta/\lambda = 0.6$ 



### Finite skin depth case, z-polarization

This is the case studied in Morgan's paper (for different profiles of grooves).



Note that  $\lambda/2\pi\delta = h/2\delta$ . In the limit  $\delta \to 0$  we previously found  $\eta_z = 1.68$ .

Take

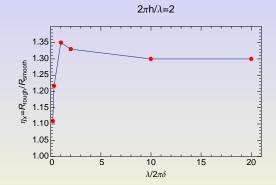
$$\boldsymbol{H} = -\hat{\boldsymbol{z}} \times \nabla \psi(\boldsymbol{x}, \boldsymbol{y})$$

At  $y \to \infty$  we have  $\psi \to H_0 y$ , and at  $y \to -\infty$  we have  $\psi \to 0$ . The equation for  $\psi$ 

$$\Delta \psi = s \frac{2i}{\delta^2} \psi$$

where the function s = 1 in the metal and zero otherwise.

#### Finite skin depth case, x-polarization



In the limit  $\delta \rightarrow 0$  we previously found  $\eta_x = 1.27$